

# MEM-284: Κυματική Διάδοση

Τμήμα Μαθηματικών και Εφαρμοσμένων Μαθηματικών, Πανεπιστήμιο Κρήτης

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$$u_n = \int \underline{f} * \underline{f}$$

## Θεώρημα αναπαράστασης

$$u_n(\mathbf{x}, t) = \int_{\Sigma_\xi} \nu_j c_{ijpq} [u_i](\xi, t) * \frac{\partial}{\partial \xi_q} G_{np}(\mathbf{x}, t; \xi) d\Sigma_\xi$$

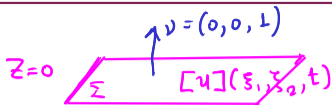
## Ισοδύναμη δύναμη

$$\boxed{e_p(\boldsymbol{\eta}, t)} = - \int_{\Sigma_\xi} \frac{\partial}{\partial \eta_q} \delta(\boldsymbol{\eta} - \boldsymbol{\xi}) \nu_j c_{ijpq} [u_i](\xi, t) d\Sigma_\xi \quad \boldsymbol{\eta} \in \Sigma$$

## Ομοιόμορφα και ιστροπικά ελαστικά μέσα

$$c_{ijpq} = \lambda \delta_{ij} \delta_{pq} + \mu (\delta_{ip} \delta_{jq} + \delta_{iq} \delta_{jp})$$

# Ισοδύναμη δύναμη - Θεώρημα αναπαράστασης



$$= u(\underline{x}, t)$$

Μετατόπιση

$$\underline{G}_{np}(x, t; \underline{\xi}) * e_p(\underline{\xi}, t) = \int_{\mathbb{R}^3} \underline{G}_{np}(x, t; \underline{\xi}, \tau) e_p(\underline{\xi}, \tau) dV_{\underline{\xi}}$$

$$u_n(\underline{x}, t) = \int_{V_0} G_{np}(\underline{x}, t; \underline{\xi}) * e_p(\underline{\xi}, t) dV_{\underline{\xi}}$$

$$\parallel$$

$$G_{np}(x, t; \underline{\xi}, \tau)$$

$$\int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty}$$

$$d\xi_1, d\xi_2$$

$$d\Sigma$$

►  $V_0$  : όγκος στον οποίο η  $e = (e_1, e_2, e_3)$  είναι μη μηδενική.

Από την άσκηση 3 της διάλεξης 10 είχαμε

$$z=0$$

$$e = \left( -\mu H(t) \delta(\eta_1) \delta(\eta_2) \frac{\partial}{\partial \eta_3} \delta(\eta_3), 0, -\mu H(t) \delta(\eta_2) \delta(\eta_3) \frac{\partial}{\partial \eta_1} \delta(\eta_1) \right)$$

$$d\tau d\Sigma_{\underline{\xi}}$$

$$u_n(\underline{x}, t) = \int_{\Sigma} \int_{\mathbb{R}^3} G_{np}(x, t; \underline{\xi}, \tau) e_p(\underline{\xi}, \tau) d\tau dV_{\underline{\xi}} = \mu \int_{\Sigma} \int_{\mathbb{R}^3} G_{n1}(x, t; \underline{\xi}, \tau) H(\tau) \delta(\xi_1) \delta(\xi_2) \frac{\partial}{\partial \xi_3} \delta(\xi_3) d\tau d\Sigma_{\underline{\xi}}$$

$$- \mu \int_{\Sigma} \int_{\mathbb{R}^3} G_{n3}(x, t; \underline{\xi}, \tau) H(\tau) \delta(\xi_2) \delta(\xi_3) \frac{\partial}{\partial \xi_1} \delta(\xi_1) d\tau d\Sigma_{\underline{\xi}}$$

$$= \mu \frac{\partial \delta(\xi_3)}{\partial \xi_3} \int_{\Sigma} \int_0^{\infty} G_{n2}(x_{\tilde{n}}, t; \tilde{\xi}, \tau) \delta(\xi_1) \delta(\xi_2) d\tau d\Sigma_{\tilde{\xi}} - \underbrace{\mu \delta(\xi_3)}_{\text{red}} \int_{\Sigma} \int_0^{\infty} G_{n3} \delta(\xi_2) \frac{\partial \delta(\xi_1)}{\partial \xi_1} d\tau d\Sigma_{\tilde{\xi}}$$

$$= -\mu \frac{\partial \delta(\xi_3)}{\partial \xi_3} \int_0^{\infty} G_{n2}(x_{\tilde{n}}, t; (0, 0, \xi_3), \tau) d\tau + \textcircled{*}$$

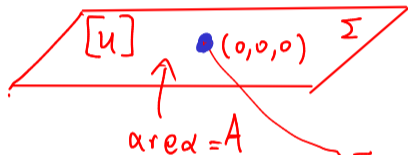
$$\textcircled{*} = -\mu \delta(\xi_3) \int_0^{\infty} d\tau \int_{\Sigma} G_{n3}(x_{\tilde{n}}, t; \tilde{\xi}, \tau) \delta(\xi_2) \left( \frac{\partial \delta(\xi_1)}{\partial \xi_1} \right) d\Sigma_{\tilde{\xi}} =$$

$$= \mu \delta(\xi_3) \int_0^{\infty} d\tau \frac{\partial}{\partial \xi_1} G_{n3}(x_{\tilde{n}}, t; (0, 0, \xi_3), \tau)$$

Προσέγγιση μη σημειακής ασυνέχειας με σημειακή πηγή.

$\mu$

$$[u_1(\xi_1, \xi_2, t)], [u_2] = [u_3] = 0, \quad (\xi_1, \xi_2) \in \Sigma$$



$$\overline{[u]}(t) = \frac{1}{A} \int_{\Sigma} [u_2](\xi_1, \xi_2, t) d\xi_1 d\xi_2$$

$$[u_2(\xi_1, \xi_2, t)] = \frac{1}{\mu} \underbrace{\mu \overline{[u]}(t) A}_{M_0(t)} \delta(\xi_1) \delta(\xi_2)$$

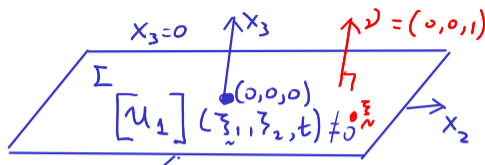
$M_0(t) \leftarrow$  ροπή της

ισοδύναμης δύναμης

$$= \frac{1}{\mu} M_0(t) \delta(\xi_1) \delta(\xi_2)$$

# Ισοδύναμη δύναμη - Θεώρημα αναπαράστασης

$u_n(\underline{x}, t) = j$   
example:



$\underline{x}$   
 $[u_2] = [u_3] = 0$

$$u_n(\underline{x}, t) = \int_{\Sigma_{\underline{\xi}}} \nu_i C_{ijpq} [u_j](\underline{\xi}, t) * \frac{\partial}{\partial \xi_q} G_{np}(\underline{x}, t; \underline{\xi}) d\Sigma_{\underline{\xi}}$$

$i=1, j=3$   
 $p=1, q=3$   
 $p=3, q=1$

$C_{13pq} = \mu, \omega \quad p=1, q=3 \quad \text{ή} \quad p=3, q=1.$

$$u_n(\underline{x}, t) = \int_{\Sigma_{\underline{\xi}}} \mu [u_2](\underline{\xi}, t) * \left\{ \frac{\partial}{\partial \xi_1} G_{n3}(\underline{x}, t; \underline{\xi}) + \frac{\partial}{\partial \xi_3} G_{n1}(\underline{x}, t; \underline{\xi}) \right\} d\Sigma_{\underline{\xi}}$$

$G_{np} = \underline{G}_{np}^P + \underline{G}_{mp}^S, \quad \underline{G}_{np}^P \xrightarrow{\theta. \text{dynam.}} u_n^P, \quad \underline{G}_{np}^S \rightarrow u_n^S$

$u_n = u_n^P + u_n^S$

$f * (g_1 + g_2) = f * g_1 + f * g_2 \quad \int g_1 + g_2 = \int g_1 + \int g_2$

# Ισοδύναμη δύναμη - Θεώρημα αναπαράστασης

Μελέτη για τα p-waves  $G_{(mp)}^p = \frac{\delta_p}{4\pi r \alpha^2} \frac{1}{r} \delta_m \delta(t - r/\alpha)$ ,  $\delta_i = \frac{x_i - \xi_i}{r}$

$$\frac{\partial \delta_i}{\partial \xi_q} = \frac{\partial}{\partial \xi_q} \left( \frac{x_i - \xi_i}{r} \right) = \frac{\frac{\partial}{\partial \xi_q} (x_i - \xi_i) r - (x_i - \xi_i) \frac{\partial}{\partial \xi_q} r}{r^2}$$

$r = \left( (x_1 - \xi_1)^2 + (x_2 - \xi_2)^2 + (x_3 - \xi_3)^2 \right)^{1/2}$

$$= \frac{-\delta_{iq} r - (x_i - \xi_i) \frac{\partial}{\partial \xi_q} r}{r^2} \quad (*)$$

$\frac{\partial}{\partial \xi_q} x_i - \frac{\partial \xi_i}{\partial \xi_q} = \delta_{iq}$

$$\frac{\partial r}{\partial \xi_q} = \frac{\partial}{\partial \xi_q} \left\{ \left( (x_1 - \xi_1)^2 + (x_2 - \xi_2)^2 + (x_3 - \xi_3)^2 \right)^{1/2} \right\} = -\frac{1}{2} r^{-1} \cdot 2(x_q - \xi_q) = -\frac{x_q - \xi_q}{r} = -\delta_q$$

$$(*) = -\frac{\delta_{iq}}{r} - \frac{1}{r} \frac{x_i - \xi_i}{r} \cdot (-\delta_q) = \frac{-\delta_{iq} + \delta_i \delta_q}{r}$$

$$\frac{\partial}{\partial \xi_q} \delta(t - r/\alpha) = \delta'(t - r/\alpha) \cdot \frac{\partial}{\partial \xi_q} (t - r/\alpha) = -\frac{1}{\alpha} \delta'(t - r/\alpha) \frac{\partial}{\partial \xi_q} r = \frac{\delta_q}{\alpha} \delta'(t - r/\alpha)$$

$$\frac{\partial \mathcal{G}_{n1}^P}{\partial \xi_3} = \frac{\gamma_1}{4\pi r \alpha^2} \frac{\partial}{\partial \xi_3} \left\{ \frac{1}{r} \gamma_n \delta(t - r/\alpha) \right\} \quad r \gg 1$$

$$\frac{\partial}{\partial \xi_3} \left\{ \frac{1}{r} \right\} \gamma_n \delta(t - r/\alpha) + \frac{1}{r} \frac{\partial}{\partial \xi_3} \gamma_n \delta(t - r/\alpha) + \frac{1}{r} \gamma_n \frac{\partial}{\partial \xi_3} \delta(t - r/\alpha)$$

$$= \frac{\gamma_3}{r^2} \gamma_n \delta(t - r/\alpha) + \frac{-\delta_{n3} + \gamma_n \gamma_3}{r^2} \delta(t - r/\alpha) + \frac{1}{r} \gamma_n \frac{\gamma_3}{\alpha} \delta'(t - r/\alpha)$$

$$\stackrel{r \gg 1}{\approx} \frac{1}{r \alpha} \gamma_n \gamma_3 \delta'(t - r/\alpha).$$

$$\frac{\partial \mathcal{G}_{n1}^P}{\partial \xi_3} \approx \frac{\gamma_n \gamma_1 \gamma_3}{4\pi r \alpha^3} \delta'(t - r/\alpha)$$

$$p=1, q=3$$

ομοίως  $\frac{\partial \mathcal{G}_{n3}^P}{\partial \xi_1} \approx \frac{\gamma_n \gamma_1 \gamma_3}{4\pi r \alpha^3} \delta'(t - r/\alpha)$

$$\left\{ \frac{\partial \mathcal{G}_{n1}^P}{\partial \xi_3} + \frac{\partial \mathcal{G}_{n3}^P}{\partial \xi_1} \right\} \approx \frac{\gamma_n \gamma_1 \gamma_3}{2\pi r \alpha^3} \delta'(t - r/\alpha)$$



$$r \gg 1$$

$$u_m^L(x, t) = \mu A \int_{\Sigma_{\tilde{x}}} [\bar{u}_L](t) \delta(\xi_1) \delta(\xi_2) * \frac{\delta_m \delta_1 \delta_3}{2\pi r \alpha^3} \delta'(t - r/\alpha) d\Sigma_{\tilde{x}} =$$

$$= \mu A \frac{\delta_m \delta_1 \delta_3}{2\pi r \alpha^3} [\bar{u}_L](t) * \delta'(t - r/\alpha)$$

$r = \|x\|$   
 $\delta_i = \frac{x_i}{\|x\|}$

$$[\bar{u}_L](t) * \delta'(t - r/\alpha) = \int_{\mathbb{R}} [\bar{u}_L](z) \delta'(t - z - r/\alpha) dz =$$

$$f(t) * g(t) = \int_{\mathbb{R}} f(z) g(t-z) dz = \int_{\mathbb{R}} g(z) f(t-z) dz = g(t) * f(t)$$

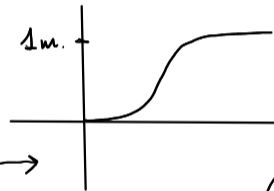
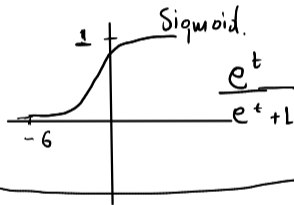
$$= [\bar{u}_L](z) \delta(t - z - r/\alpha) \Big|_{-\infty}^{+\infty} - \int_{\mathbb{R}} [\bar{u}_L]'(z) \delta(t - z - r/\alpha) dz$$

$$= - \int_{\mathbb{R}} [\overline{u_1}]'(\tau) \delta(t - \tau - r/\alpha) d\tau = - [\overline{u_1}]'(t - r/\alpha)$$

$$u_{\eta}^P(\underline{x}, t) = -\mu A \frac{\delta_{\eta} \delta_1 \delta_3}{2\pi\rho r \alpha^3} [\overline{u_1}]'(t - r/\alpha), \quad \eta = 1, 2, 3$$

Παράδειγμα

$$[\overline{u_1}](t) = \frac{e^{t-6}}{e^{t-6} + 1}$$



$$[\overline{u_1}]'(t) = \frac{e^{t-6} (e^{t-6} + 1) - e^{t-6} e^{t-6}}{(e^{t-6} + 1)^2} = \frac{e^{t-6}}{(e^{t-6} + 1)^2}$$



$$u_{\eta}^P(x_{\sim}, t) = -\mu A \frac{\gamma_1 \gamma_2 \gamma_3}{4\pi r d^3} \frac{e^{t-6-r/\alpha}}{(e^{t-6-r/\alpha} + 1)^2}$$

Παράδειγμα:

$$\mu = \rho b^2$$



θ.δ.ο

$$[\bar{u}](t) = H(t), t \in \mathbb{R}$$

$$[\bar{u}]'(t) = \delta(t)$$

$$H'(t) = \delta(t)$$

$$H(t) = \int_{-\infty}^t H'(z) dz = \begin{cases} 1, & t > 0 \\ 0, & t < 0 \end{cases}$$

$$\int_{-\infty}^t \delta(z) dz = H(t)$$



$$u_{\eta}^P(x_{\sim}, t) \propto \delta(t - r/\alpha)$$

# Ισοδύναμη δύναμη - Θεώρημα αναπαράστασης

$$\gamma_i = \frac{x_i}{r}$$

$$r = (x_1^2 + x_2^2 + x_3^2)^{1/2}$$

